# An effect-theoretic reconstruction of quantum theory 

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## Why Quantum Theory?

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Its mathematical description is not particularly compelling:

- Systems are described by C*-algebras.
- States are density matrices.
- Dynamics are completely positive maps.
- Measurement outcomes are governed by the trace rule.
- Composite systems are formed using the tensor product.


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Not clear at all why this describes nature so well.

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Find sensible physical requirements from which it follows.

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Find sensible physical requirements from which it follows.

If successful, we can say:
Quantum theory describes nature because "it couldn't have been any other way"
(without nature being that much weirder)

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In this talk:
"Any theory with well-behaved pure maps is quantum theory"
All axioms taken from effectus theory

## A suitable framework

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An effectus $\approx$ 'generalised generalised probabilistic theory'
real numbers $\Rightarrow$ effect monoids
vector spaces $\Rightarrow$ effect algebras.

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2. The maps $v, w:(I+I)+I \rightarrow I+I$ given by

$$
v=\left[\left[\kappa_{1}, \kappa_{2}\right], \kappa_{2}\right] \text { and } w=\left[\left[\kappa_{2}, \kappa_{1}\right], \kappa_{2}\right] \text { are jointly monic }
$$

(i.e. $v \circ f=v \circ g$ and $w \circ f=w \circ g$, then $f=g$ ).

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- Sets (or more generally any topos).


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- Opposite of category of order unit spaces In particular any (causal) general probabilistic theory.
- Opposite category of von Neumann algebras


## Basic definitions and consequences

- Partial maps: $f: X \rightarrow Y+I$.
- States: $\operatorname{St}(X):=\operatorname{Hom}(I, X)$.
- Effects: $\operatorname{Eff}(X):=\operatorname{Hom}(X, I+I)$.
- Scalars: $\operatorname{Hom}(I, I+I)$.


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- The states form an abstract convex set.
- The effects form an effect algebra.
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- The states form an abstract convex set.
- The effects form an effect algebra.
- The partial maps preserve this structure.

Definition of effectus is basically chosen to make these things true

## Effect algebras

## Definition

An effect algebra $\left(E, 0,1,+,(\cdot)^{\perp}\right)$ is a set $E$ with partial commutative associate "addition" + and involution $(\cdot)^{\perp}$ such that

- $\left(x^{\perp}\right)^{\perp}=x$,
- $x+x^{\perp}=1$,
- If $x+1$ is defined, then $x=0$.


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- In particular: set of effects of C*-algebra.


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- Any Boolean algebra
- Any interval $[0, u]$ with $u \geqslant 0$ in an ordered vector space
- In particular: set of effects of $C^{*}$-algebra.

Note: Effect algebra is partially ordered by $x \leqslant y$ iff $\exists z: x+z=y$.

## Baby effectus

Definition
A Effect theory is a category $\mathbf{B}$ with designated object $/$ such that $\operatorname{Hom}(A, I)$ is an effect algebra, and for any $f: B \rightarrow A$ :
$0 \circ f=0,(p+q) \circ f=(p \circ f)+(q \circ f)$.

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$0 \circ f=0,(p+q) \circ f=(p \circ f)+(q \circ f)$.
Very basic structure, we need more assumptions!

## Compressions and filters

A compression for $q: A \rightarrow I$ is a map $\pi_{q}:\{A \mid q\} \rightarrow A$ with $1 \circ \pi_{q}=q \circ \pi_{q}$,

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A filter for $q: A \rightarrow I$ is a $\operatorname{map} \xi_{q}: A \rightarrow A_{q}$ with $1 \circ \xi \leqslant q$,

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## Quotient and Comprehension: All the adjunctions!

$$
\underset{\substack{\text { Quotient } \\(X, p) \mapsto X / p}}{\operatorname{Pred}_{\square}(\mathbf{C})}
$$

$$
\begin{gathered}
\operatorname{Pred}_{\square}(\mathbf{C}): \\
\text { Objects are }(X, p: X \rightarrow I) . \\
\text { Morphisms: } f:(X, p) \rightarrow(Y, q) \text { is } \\
f: X \rightarrow Y \text { with } p^{\perp} \geqslant q^{\perp} \circ f .
\end{gathered}
$$

Source: arXiv:1512.05813, p. 97

See also: Cho, Jacobs, Westerbaan ${ }^{2}$ 2015. Quotient-Comprehension Chains

## Example

Let $\mathrm{Mat}_{\mathbb{C}}^{\text {op }}$ be the opposite category of positive sub-unital maps $f: M_{n}(\mathbb{C}) \rightarrow M_{m}(\mathbb{C})$. I.e $a \geqslant 0 \Longrightarrow f(a) \geqslant 0$ and $f(1) \leqslant 1$.

## Example

Let $\mathbf{M a t} \mathbf{t}_{\mathbb{C}}^{\text {op }}$ be the opposite category of positive sub-unital maps $f: M_{n}(\mathbb{C}) \rightarrow M_{m}(\mathbb{C})$. I.e $a \geqslant 0 \Longrightarrow f(a) \geqslant 0$ and $f(1) \leqslant 1$.

An effect then corresponds to $q \in M_{n}(\mathbb{C})$ with $0 \leqslant q \leqslant 1$.
Write $q=\sum_{i} \lambda_{i} q_{i}$ with $\lambda_{i}>0, q_{i} q_{j}=\delta_{i j} q_{i}$.
Define $\lceil q\rceil=\sum_{i} q_{i} .\lfloor q\rfloor=\sum_{i ; \lambda_{i}=1} q_{i}$.

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Define $\lceil q\rceil=\sum_{i} q_{i} .\lfloor q\rfloor=\sum_{i ; \lambda_{i}=1} q_{i}$.
The projection $\pi_{q}: M_{n}(\mathbb{C}) \rightarrow\lfloor q\rfloor M_{n}(\mathbb{C})\lfloor q\rfloor$ is a compression.
$\xi_{q}:\lceil q\rceil M_{n}(\mathbb{C})\lceil q\rceil \rightarrow M_{n}(\mathbb{C})$ with $\xi_{q}(p)=\sqrt{q} p \sqrt{q}$ is a filter.

## Images, kernels and cokernels

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An effect theory has images, and for all sharp effects compressions and filters if and only if the category has all kernels and cokernels.

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In fact: compressions are kernels, and filters for sharp effects are cokernels.
$\Rightarrow$ filters are "fuzzy" cokernels.

## Pure maps

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Motivation: In Mat ${ }_{\mathbb{C}}^{\text {op }}$ a map $f: M_{n}(\mathbb{C}) \rightarrow M_{m}(\mathbb{C})$ is pure iff $\exists V: \mathbb{C}^{n} \rightarrow \mathbb{C}^{m}$ such that $f(a)=V a V^{\dagger}$ for all a.

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## Remark

From definition it is not clear that pure maps are closed under composition. But: In Mat ${ }_{\mathbb{C}}^{\mathrm{op}}$ it is true.
Also: there is an obvious dagger on pure maps in Mat ${ }_{\mathbb{C}}^{\mathrm{Op}}$.

## Pure effect Theories

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5. If $\pi_{q}$ is a compression for sharp $q$, then $\pi_{q}^{\dagger}$ is a filter for $q$.
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## PET examples

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- vNA ${ }_{\text {ncpsu }}^{\circ \mathrm{p}}$ : von Neumann algebras with normal completely positive sub-unital maps between them.
- Category of real C*-algebras.
- EJA psu: positive sub-unital maps between Euclidean Jordan algebras.


## Euclidean Jordan algebras

## Definition

A Euclidean Jordan algebra (EJA) $(E,\langle\cdot, \cdot\rangle, *, 1)$ is a real Hilbert space with a product that satisfies $\forall a, b, c$ :

$$
a * 1=a \quad a * b=b * a \quad a *\left(b * a^{2}\right)=(a * b) * a^{2} \quad\langle a * b, c\rangle=\langle b, a * c\rangle
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We have an order $a \geqslant 0 \Longleftrightarrow \exists b: a=b * b:=b^{2}$.
Example: $M_{n}(F)^{\text {sa }}$ - self-adjoint matrices over $F=\mathbb{R}, \mathbb{C}, \mathbb{H}$ with $A * B:=\frac{1}{2}(A B+B A)$ and $\langle A, B\rangle:=\operatorname{tr}(A B)$.

## HTA



Me explaining why Jordan algebras are cool:


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- If $\operatorname{Eff}(A) \cong[0,1]$ then $A \cong 1$.

Operational effect theory $\approx$ generalized probabilistic theory

## Main result 1: Everything is a Jordan algebra

Theorem
Let $\mathbf{B}$ be an operational PET. Then there is a functor $F: \mathbf{B} \rightarrow \mathbf{E J A}_{p s u}^{\circ \mathrm{op}}$ with $F(\operatorname{Eff}(A)) \cong \operatorname{Eff}(F(A))$.

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(If $\forall p: p \circ f=p \circ g$ then $f=g$ )
"Operational PETs are Euclidean Jordan algebras"

## Monoidal effect theories

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How to go from Jordan algebras to quantum theory?
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## Definition

An effect theory is monoidal when it is monoidal with I as unit such that tensor preserves addition. A PET is monoidal if the subcategory of pure maps is in addition also monoidal.

## Quantum Theory Reconstructed

Theorem
Let $\mathbf{B}$ be a monoidal operational PET. Then there is a functor $F: \mathbf{B} \rightarrow \mathbf{C}^{\text {op }}$ with $F(\operatorname{Eff}(A)) \cong \operatorname{Eff}(F(A))$ where $\mathbf{C}$ is the category of real or complex $C^{*}$-algebras.

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Furthermore, if effects separate maps, then it is faithful and C*-algebras must be complex.

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Furthermore, if effects separate maps, then it is faithful and C*-algebras must be complex.

Recall the assumptions:

1. All maps have images.
2. When $q$ is sharp, $q^{\perp}$ is sharp.
3. All effects have filters and compressions.
4. The pure maps form a monoidal dagger-category.
5. If $\pi_{q}$ is a compression for sharp $q$, then $\pi_{q}^{\dagger}$ is a filter for $q$.
6. Compressions for sharp $q$ are isometries: $\pi_{q}^{\dagger} \circ \pi_{q}=$ id.

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Future work:

- Minimality of conditions?
- How much can be done in abstract setting?
- Can we get Jordan algebras over different fields?
- Characterize infinite-dimensional quantum theory?


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Dagger and dilations in the category of von Neumann algebras
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Thank you for your attention

